

# QCD thermodynamics with $N_f=3,2+1$ near the continuum limit at realistic quark mass

## ~ status report ~

Takashi Umeda (BNL)

### *BNL*

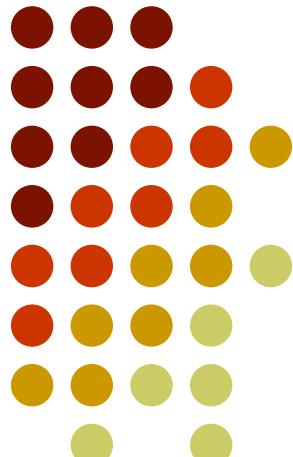
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Peter Petreczky

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Norman Christ  
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Olaf Kaczmarek  
Edwin Laermann  
Chuan Miao  
Stanislav Shcheredin  
Jan van der Heide  
Sonke Wissel





# Motivation

- The critical parameters of the QCD transition and EoS from first principle calculation (Lattice QCD)  
 $T_c$ ,  $\varepsilon_c$ , phase diagram, small  $\mu_B$ , etc...
- These are very important for Heavy Ion Phenomenology  
many phenomenological models  
based on the parameters from lattice QCD results

*More accurate determination of these params. is required !*



from recent studies  
these results strongly depend on quark mass &  $N_f$

Our aim is thermodynamics at almost realistic quark mass &  $N_f$   
(2+1)-flavor with, pion mass  $\sim 200\text{MeV}$ , kaon mass  $\sim 500\text{MeV}$

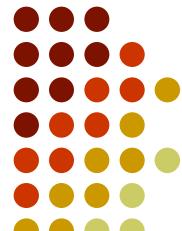


# Our Strategy

For "the almost realistic quark mass at  $N_f=2+1$ "

- pion mass  $\sim 200\text{MeV}$ , kaon mass  $\sim 500\text{MeV}$

- Choice of quark action
  - Staggered type quark action
- huge computational resource is required
  - QCDOC machine, APE-Next machine
- continuum limit
  - $N_t=4,6(,8)$  →  $a=0.2,0.16(,0.1)\text{fm}$
  - improved action for reliable continuum limit with not so fine lattices



# Our Computers

US/RBRC QCDOC

20.000.000.000.000 ops/sec



- critical temperature
- equation of state
- hadron properties in matter

BI – apeNEXT

5.000.000.000.000 ops/sec



today: 1.6 TFlops

<http://www.quark.phy.bnl/~hotqcd>



# Lattice Action

## improved Staggered action : p4-action

*Karsch, Heller, Sturm (1999)*

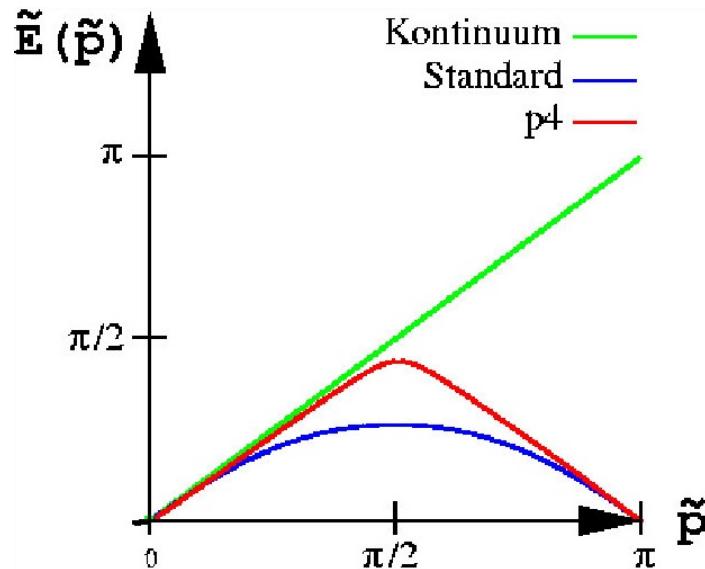
- gluonic part : Symanzik improvement scheme
  - remove cut-off effects of  $O(a^2)$
  - tree level improvement  $O(g^0)$
- fermion part : improved staggered fermion
  - remove cut-off effects & improve rotational sym.
  - improve flavor symmetry by smeared 1-link term

$$S_F(N_\tau, N_\sigma) = \sum_{n, \dot{n}} \sum_{\mu} \eta(n_\mu) \bar{\chi}_n \left( \frac{3}{8} \left[ \frac{1}{1+6\omega} \left( \begin{array}{c} \leftarrow \circ \rightarrow \\ \text{--- --- ---} \end{array} \right) + \omega \sum_{\nu \neq \mu} \begin{array}{c} \leftarrow \uparrow \downarrow \rightarrow \\ \text{--- --- ---} \\ \leftarrow \uparrow \downarrow \rightarrow \\ \text{--- --- ---} \end{array} \right] \right. \right. \\ \left. \left. + \frac{1}{48} \sum_{\nu \neq \mu} \left[ \begin{array}{cccc} \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow \\ \text{--- --- ---} & \text{--- --- ---} & \text{--- --- ---} & \text{--- --- ---} \\ \downarrow & \downarrow & \downarrow & \downarrow \\ \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow & \leftarrow \uparrow \downarrow \rightarrow \\ \text{--- --- ---} & \text{--- --- ---} & \text{--- --- ---} & \text{--- --- ---} \end{array} \right] \right] \right) \chi_{n'} + m_q \sum_n \bar{\chi}_n \chi_n \right)$$



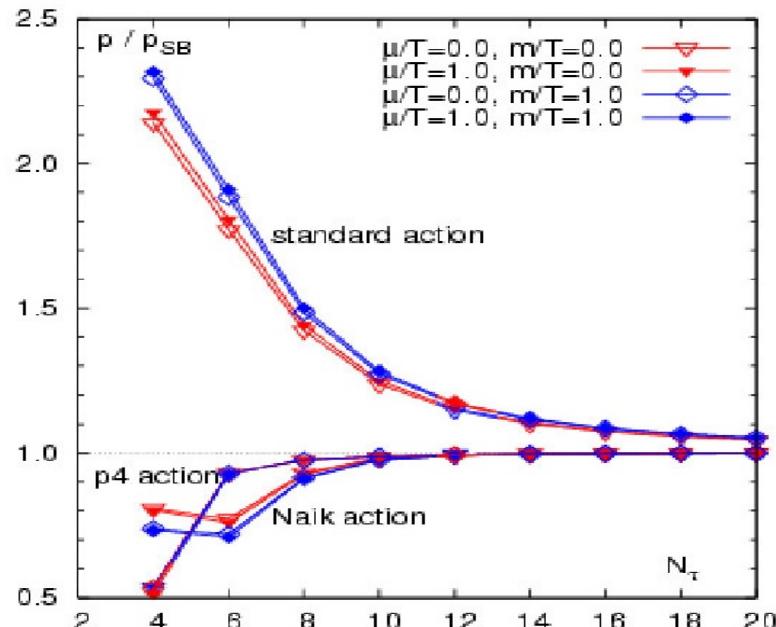
# Properties of the p4-action

## Dispersion relation



The free quark propagator is rotational invariant up to  $O(p^4)$

## pressure in high T limit



Bulk thermodynamic quantities show drastically reduced cut-off effects

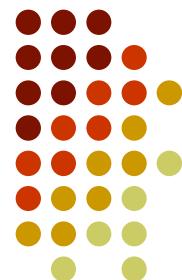
flavor sym. is also improved by fat link



## Numerical results

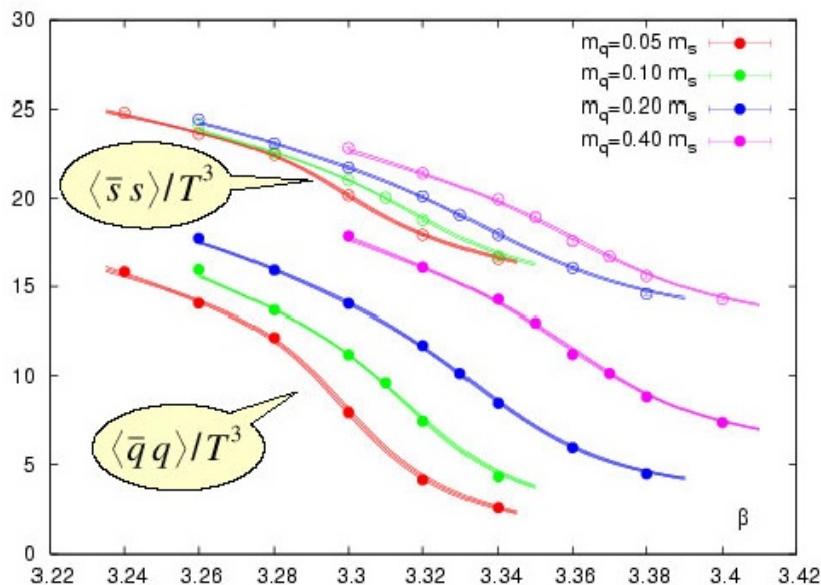
Calculation for Critical temperature

# Order parameters



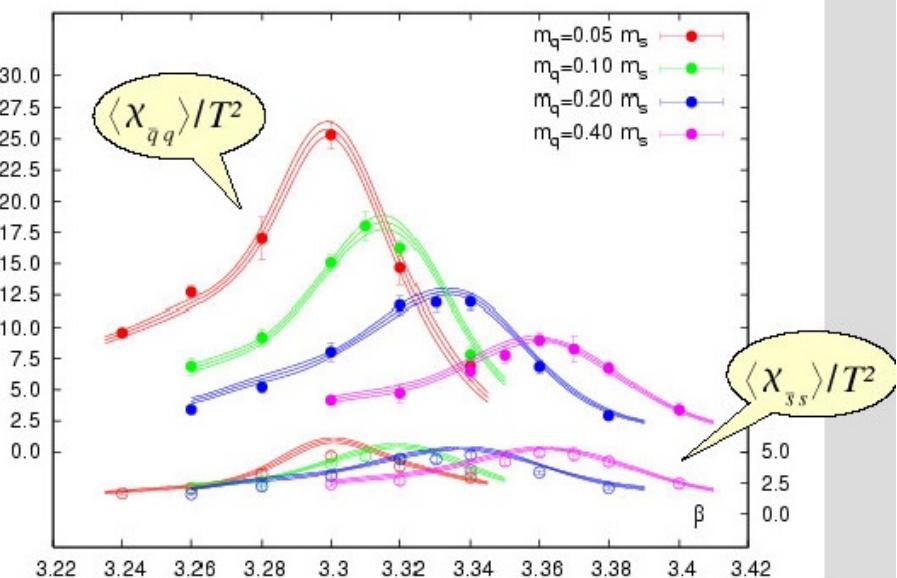
(2+1)-flavor,  $8^3 \times 4$  lattice

strange- & light quark-chiral condensate:



strange- & light quark-chiral susceptibility:

$$\langle \chi_{\bar{q}q} \rangle \equiv \langle (\bar{q}q)^2 \rangle - \langle \bar{q}q \rangle^2$$



- multi-histogram method (Ferrenberg-Swendsen) is used
- Transition becomes stronger for smaller light quark masses
- $\beta_c$  are determined by peak positions of the susceptibilities

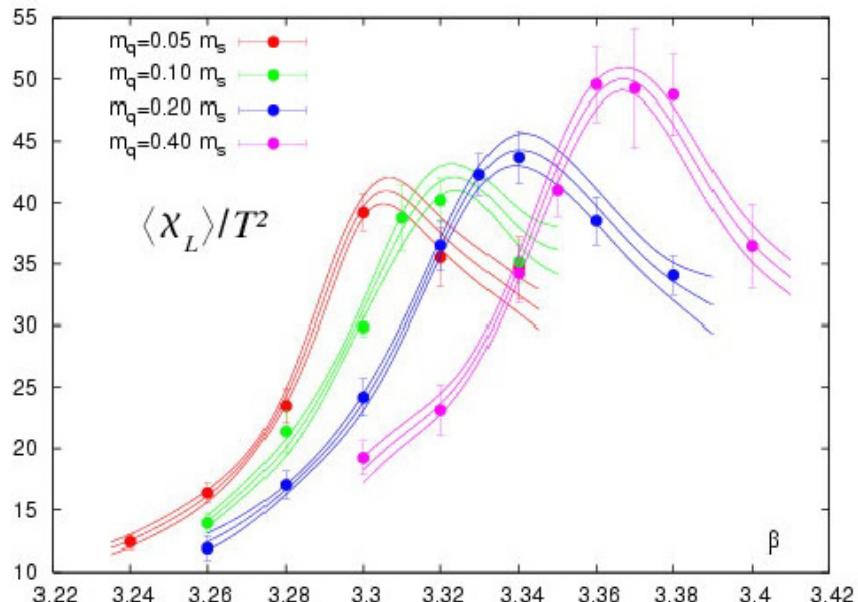
# Order parameters



(2+1)-flavor,  $8^3 \times 4$  lattice

Polyakov loop susceptibility:

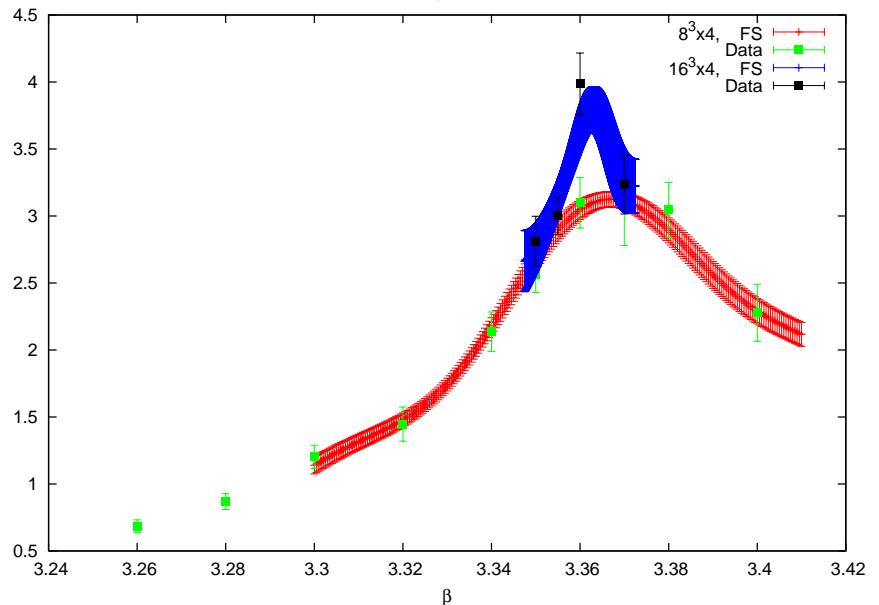
$$\langle \chi_L \rangle \equiv \langle L^2 \rangle - \langle L \rangle^2$$



$16^3 \times 4$  &  $8^3 \times 4$  lattices

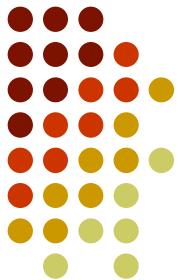
Polyakov loop susceptibility:

$$\langle \chi_L \rangle \equiv \langle L^2 \rangle - \langle L \rangle^2$$



- $\beta_c$  are determined by peak positions of the susceptibilities  
→ consistent with  $\beta_c$  from chiral susceptibility
- Transition becomes stronger for larger volume

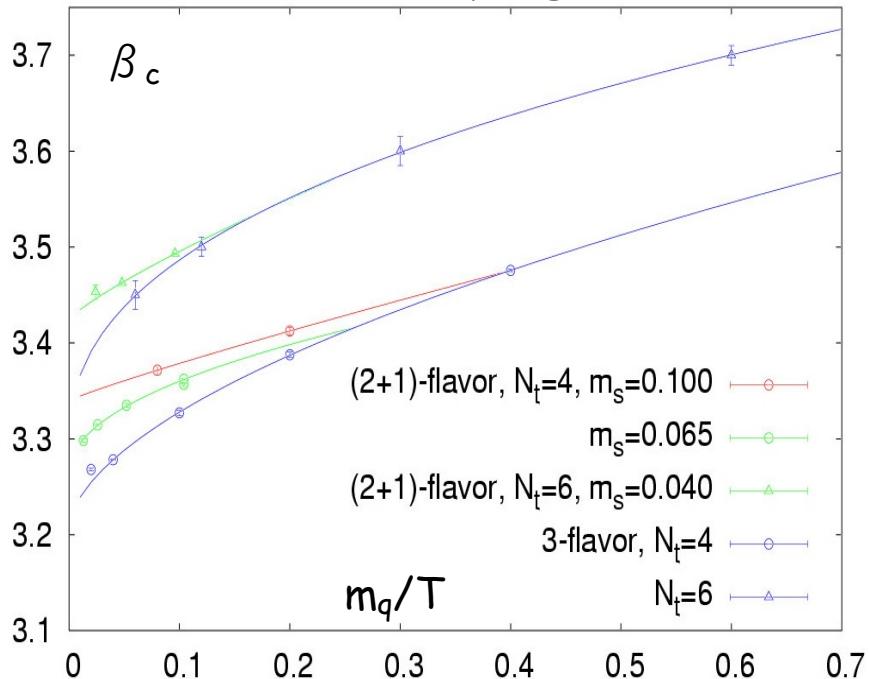
# Critical temperature



## (1) critical beta search

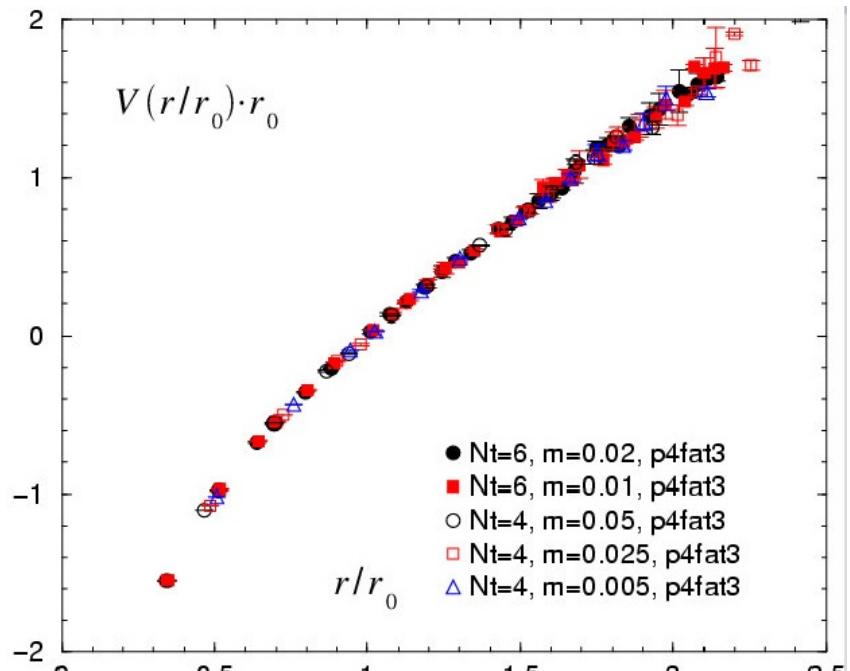
- from the chiral susceptibilities
- fits with power laws

Critical couplings:



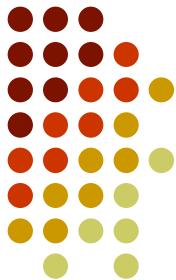
## (2) scale determination

from static quark potential  
Sommer scale & string tension



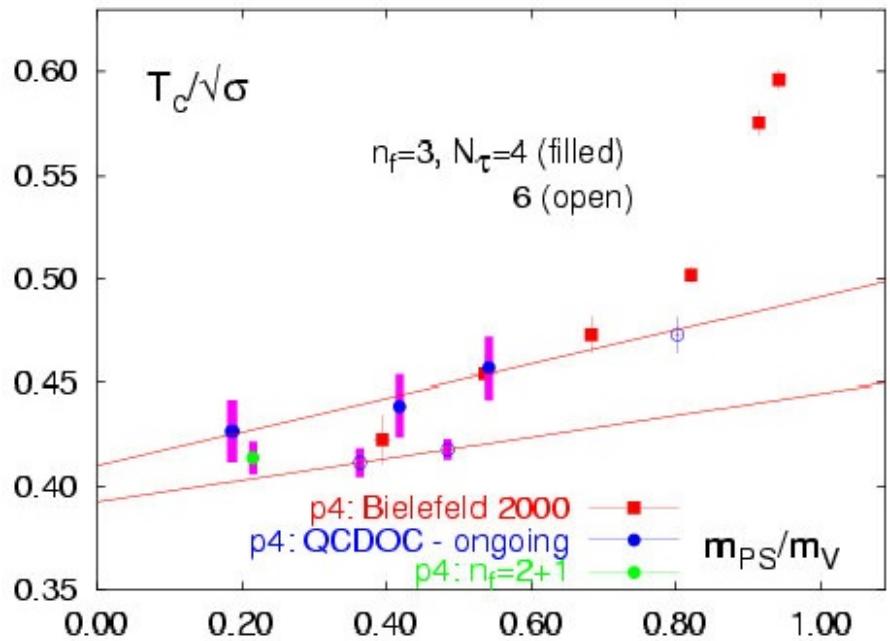
Almost no mass & cutoff dep.  
in potential scaled by  $r_0$

# Critical temperature

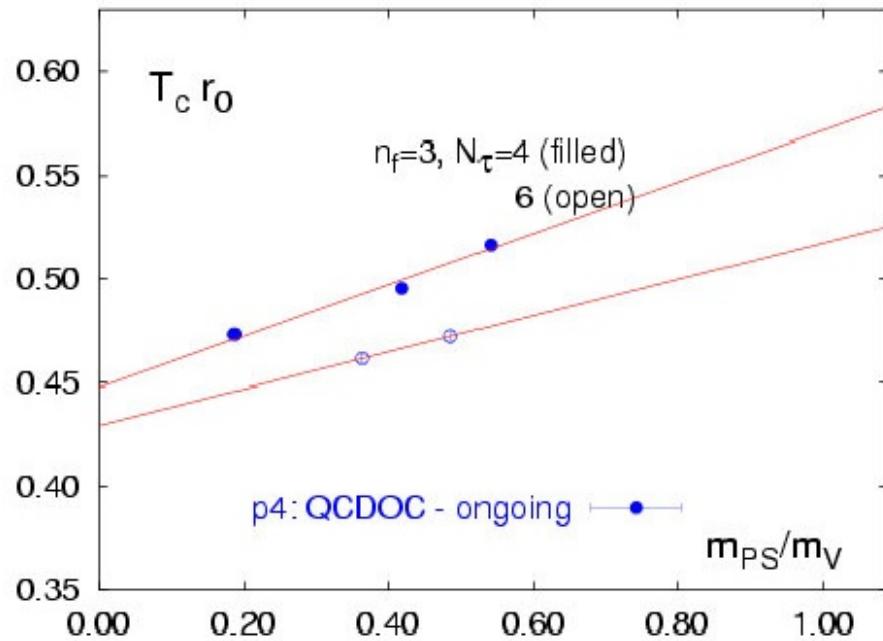


3-flavor results

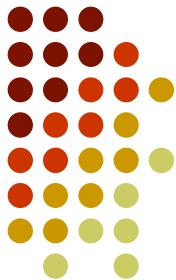
from string tension:



from Sommer scale:



- The cut-off effect in  $T_c$  is about 5% in  $m_q=0$  limit of 3-flavor QCD
- Results is consistent with previous Bielefeld result



## Summary

Critical coupling, temperature

- 3-flavor QCD
  - $m_{\text{pi}}/m_{\text{rho}} \geq 0.2$ ,  $N_\sigma=8,16,32$ ,  $N_\tau=4,6$
- (2+1)-flavor QCD
  - $m_q/m_s \geq 0.05$ ,  $N_\sigma=8,16$ ,  $N_\tau=4$

## Outlook

(2+1)-flavor  $N_\tau=4,6$

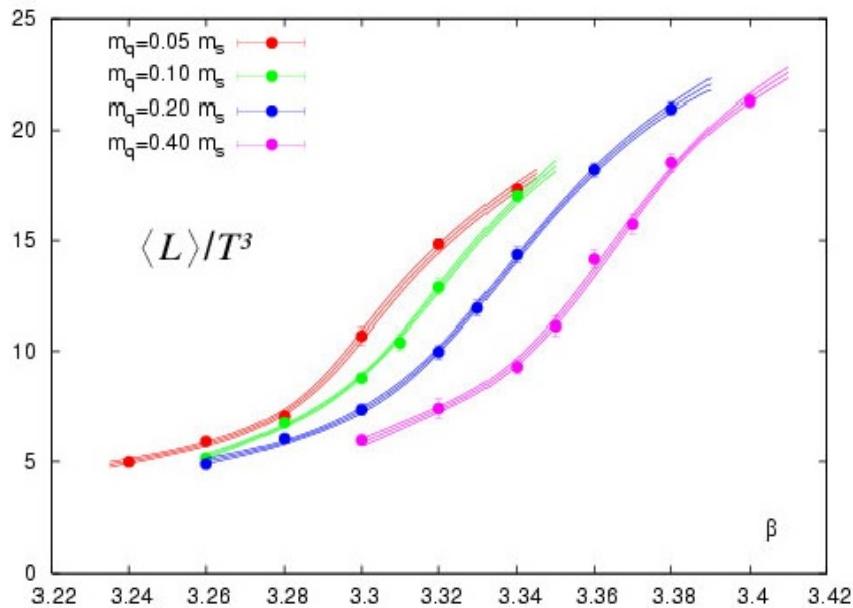
- determination of  $T_c$
- Calculation of EoS
- etc...



# Order parameters

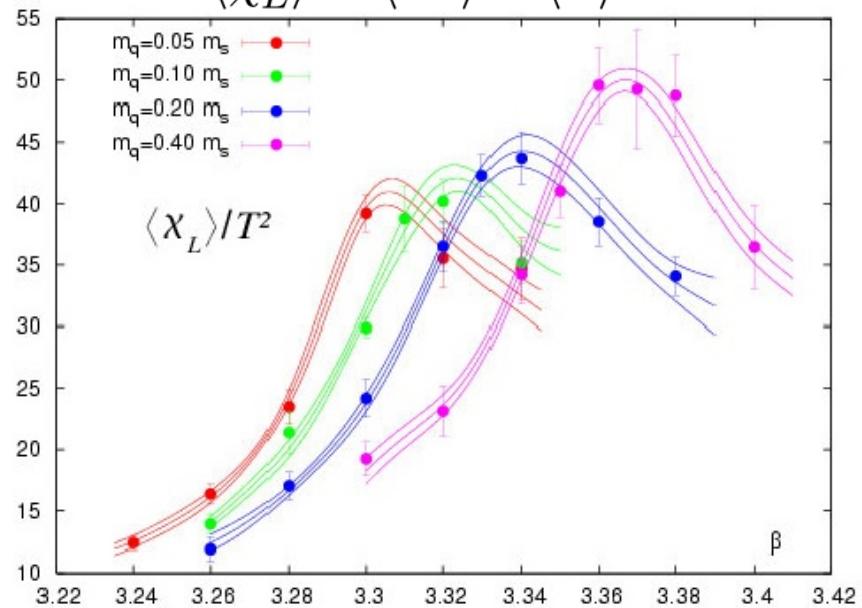
(2+1)-flavor,  $8^3 \times 4$  lattice

Polyakov loop:



Polyakov loop susceptibility:

$$\langle \chi_L \rangle \equiv \langle L^2 \rangle - \langle L \rangle^2$$



- $\beta_c$  are determined by peak positions of the susceptibilities  
→ consistent with  $\beta_c$  from chiral susceptibility

(2)finite V fig.



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